MODELING RADIAL VELOCITIES AND ECLIPSE PHOTOMETRY OF THE KEPLER TARGET KIC 4054905: AN OSCILLATING RED GIANT IN AN ECLIPSING BINARY

M. Benbakoura[1,2], P. Gaulme[3,4,5], J. McKeever[4,6], P. G. Beck[7,8], J. Jackiewicz[4] and R. A. García[1,2]

Abstract. Asteroseismology is a powerful tool to measure the fundamental properties of stars and probe their interiors. This is particularly efficient for red giants because their modes are well detectable and give information on their deep layers. However, the seismic relations used to infer the mass and radius of a star have been calibrated on the Sun. Therefore, it is crucial to assess their accuracy for red giants which are not perfectly homologous to it. We study eclipsing binaries with a giant component to test their validity. We identified 16 systems for which we intend to compare the dynamical masses and radii obtained by combined photometry and spectroscopy to the values obtained from asteroseismology. In the present work, we illustrate our approach on a system from our sample.

Keywords: asteroseismology, methods: data analysis, binaries: eclipsing, stars: evolution, stars: oscillations

1 Introduction

Over the last 10 years, asteroseismology has been the most efficient tool to probe stellar interiors. In this field, red giants are of particular interest because, on the one hand, they are solar-like pulsators, i. e., their oscillations are stochastically excited by granulation, and on the other hand, the presence of mixed-modes in their power spectrum allows to probe their core (e.g. Beck et al. 2011, 2012; Mosser et al. 2011). Moreover, the huge sample of observed red giants from *Kepler* (Borucki et al. 2010) and CoRoT (Baglin et al. 2006), including more than 30,000 stars, allows to make statistical inferences as a function of different parameters.

Obtaining masses and radii of red giants allows to know more about the fate of the Sun. These parameters can be directly obtained by measuring the effective temperature as ell as the global seismic parameters $\Delta\nu$ and ν_{max}, by applying the so-called global seismic scaling relations (Brown et al. 1991; Kjeldsen & Bedding 1995):

$$\frac{R_*}{R_\odot} = \left(\frac{\Delta\nu_\odot}{\Delta\nu}\right)^2 \left(\frac{\nu_{\mathrm{max}}}{\nu_{\mathrm{max},\odot}}\right) \left(\frac{T_{\mathrm{eff}}}{T_{\mathrm{eff},\odot}}\right)^{1/2} \tag{1.1}$$

$$\frac{M_*}{M_\odot} = \left(\frac{\Delta\nu_\odot}{\Delta\nu}\right)^4 \left(\frac{\nu_{\mathrm{max}}}{\nu_{\mathrm{max},\odot}}\right)^4 \left(\frac{T_{\mathrm{eff}}}{T_{\mathrm{eff},\odot}}\right)^{3/2} \tag{1.2}$$

These relations have been calibrated on the Sun and are based on the assumption that solar-like pulsators are homologous to our star. However, in red giants the aspect ratio (radius of the radiative zone divided by the radius of the star) may be smaller than in main-sequence stars by more than one order of magnitude. Moreover, the envelope of the former is significantly less dense because of the expansion. The departure from homology and asymptotic regime might introduce a systematic error in the seismic relations (e. g. Mosser et al. 2013;

[1] IRFU, CEA, Université Paris-Saclay, F-91191 Gif-sur-Yvette, France

[2] Université Paris Diderot, AIM, Sorbonne Paris Cité, CEA, CNRS, F-91191 Gif-sur-Yvette, France

[3] Max-Planck-Institut für Sonnensystemforschung, Justus-von-Liebig-Weg 3, 37077, Göttingen, Germany

[4] Department of Astronomy, New Mexico State University, P.O. Box 30001, MSC 4500, Las Cruces, NM 88003-8001, USA

[5] Physics Department, New Mexico Institute of Mining and Technology, 801 Leroy Place, Socorro, NM 87801, USA

[6] Department of Astronomy, Yale University, 52 Hillhouse Avenue, New Haven, CT 06511, USA

[7] Instituto de Astrofísica de Canarias, E-38200, La Laguna, Tenerife, Spain

[8] Universidad de La Laguna, Dpto. de Astrofísica, E-38205, La Laguna, Tenerife, Spain

Rodrigues et al. 2017, and references therein) Therefore, the scaling relations should be studied carefully and it is crucial to test their accuracy.

One way of testing the validity of seismic scaling relations is to compare the values of stellar mass and radii from seismology with results from independent methods. Gaulme et al. (2016) and Rawls et al. (2016) considered eclipsing binary systems with at least one giant component. For such systems, the mass and radius of both components can be accurately measured by combining photometry and spectroscopy. Indeed, the simultaneous modeling of the eclipses and the radial velocities allows to retrieve, on the one hand, the orbital parameters of the systems such as eccentricity, inclination, longitude of periastron, and semi-major axis, and on the other hand, global properties of its components. On a sample of 10 systems, Gaulme et al. (2016) compared the values of masses and radii obtained through this method (the dynamical values) with the seismic inferences. They found that the latter were systematically overestimated by $\sim 5\%$ for the radius and $\sim 15\%$ for the mass.

To investigate further this result, we aim at extending this set of stars to a larger fraction of multiple systems containing at least one pulsating red giant. Hence, we have identified 16 new binary systems with a giant component. All of them have been observed by *Kepler*. We combine the photometric light curves with spectroscopic follow-up observations from the ARCES spectrometer on the 3.5m telescope at Apache Point Observatory (APO), New Mexico, and the HERMES spectrometer (Raskin et al. 2011) on the 1.2m Mercator Telescope in La Palma to measure the dynamical masses and radii and compare them to the seismic inferences. The latter were obtained from *Kepler* data prepared with KADACS software (García et al. 2011), which was analyzed by the asteroseimic pipeline A2Z (Mathur et al. 2010).

In this work, we show our preliminary analysis of KIC 4054905. Section 2 details how we measured the dynamical mass and radius of the giant. In section 3, we explain how we prepared the light curve from which the global seismic parameters were computed. Then, we show the result of the comparison between seismic estimates and dynamical measurements. Finally, in section 4, we present the perspectives of our work and the suggestions to improve the scaling relations.

2 Measuring the dynamical masses and radii

The photometric light curves were obtained from the Target Pixel Files, which we integrated to produce time series. They were then flattened so that only the eclipses remain. We modeled the orbit with the JKTEBOP code (Southworth 2013), which uses a Markov chain Monte Carlo (MCMC) method to fit the sum of the fractional radii $\frac{R_1+R_2}{a}$, where R_1, R_2 and a are the radius of the giant, that of the companion and the semi-major axis, respectively, the ratio of the radii $\frac{R_2}{R_1}$, the orbital inclination and eccentricity, the longitude of periastron, the brightness ratio $\left(\frac{T_{\mathrm{eff},2}}{T_{\mathrm{eff},1}}\right)^4$, orbital period and reference time for the primary eclipse. Concerning the eclipses, we used the same convention as Gaulme et al. (2016), i.e., the primary eclipse refers to the companion eclipsing the giant. We set its orbital phase to 0. For the limb darkening of the companion, we assumed a linear law. We did not fit it because of the small influence it has compared to that of the giant and the resulting difficulty to converge on this parameter. Concerning the giant, we assumed a quadratic law and computed guesses with the JKTLD routine, which uses limb-darkening tables to interpolate them in the *Kepler* bandpass. We fixed the order-2 coefficient and fitted the linear one.

We combined the modeling of the eclipses with a two-year spectroscopic follow-up done mostly at APO. We used the Echelle spectrometer ARCES located at the 3.5m telescope. The data was reduced as described by Rawls et al. (2016). We used the broadening function (BF) technique as outlined by Rucinski (2002) to extract the radial velocities. This method relies on solving a convolution equation and requires a template spectrum to compare the positions of the spectral lines. In this work, we used templates from the PHOENIX stellar atmosphere model Husser et al. (2013). Since the giant is in general three to ten times more luminous than its companion, its lines are significantly easier to detect. This is why we used a template of a main-sequence star to compute the BF. Its temperature was chosen the closest to that of the companion. Combining radial velocities to the eclipses allows to break the degeneracy in the model. Thus, it is possible to obtain the semi-major axis of the system and the separate masses and radii of each component. Figure 1 illustrates the modeling of combined eclipses and radial velocities of KIC 4054905. For this star, the dynamical analysis gives $M_{\mathrm{giant}} = 1.05\ M_{\odot}$ and $R_{\mathrm{giant}} = 8.24\ R_{\odot}$.
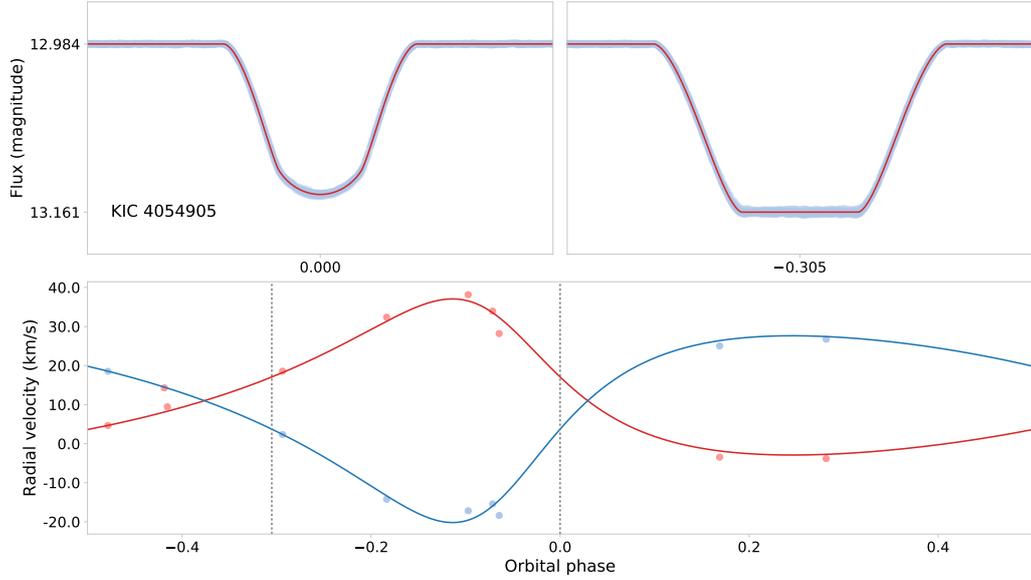
Fig. 1. Eclipses and radial velocities of KIC 4054905. *Top left panel:* Zoom on the primary eclipse (companion eclipsing the giant. The blue dots represent the processed phase-folded *Kepler* light curve and the red line, the fit of JKTEBOP. *Top right panel:* zoom on the secondary eclipse (giant eclipsing the companion). *Bottom panel:* Radial velocities of the components of the systems. The red and blue dots represent the measurements done for the giant and the companion at APO, respectively. The red and blue solid lines represent the fit of JKTEBOP for each component, respectively. The vertical dotted lines the orbital phases corresponding to the primary and second eclipse.

3 Seismic inferences

We prepared the light curves from the *Kepler* Target Pixel Files as described by García et al. (2011). In this procedure, we inserted an intermediate step to remove the eclipses from the signal. To that end, we used the Eclipsing Binary Catalog, which provides information on eclipses timing and durations (Abdul-Masih et al. 2016). Removing the eclipses in the time series generates gaps in the data. To eliminate their signature, we inpainted the light curves (García et al. 2014; Pires et al. 2015). From the obtained time series, we computed the power spectra and used the pipeline A2Z to compute the global seismic parameters $\Delta\nu$ and ν_{max}. For KIC 4054905, we obtained $M_{\mathrm{seismo}} = 1.3 \pm 0.1\ M_\odot$ and $R_{\mathrm{seismo}} = 9.3 \pm 0.7\ R_\odot$.

Figure 2 presents the comparison between dynamical measurements and seismic estimates of masses and radii for the sample of Gaulme et al. (2016) and KIC 4054905. Our preliminary result is in good agreement with their study since both mass and radius of our star are over estimated by the scaling relations.

4 Conclusion

Our preliminary result is in agreement with those of Gaulme et al. (2016), suggesting that the scaling relations systematically overestimate mass and radius. However, more spectroscopic observations of the other targets of our sample are needed to draw more robust conclusions. Extending the sample is crucial to test their accuracy. The departure of red giants from the asymptotic regime was already discussed by Mosser et al. (2013). Recently, Rodrigues et al. (2017) proposed to improve the scaling relations by taking into account that red giants depart from homology to the Sun. Other reasons for discrepancy between dynamical measurements and seismic estimates should also be investigated, such as the effect of other observables. For instance, Corsaro et al. (2017) found of metallicity influencing granulation in stellar clusters. Thus, considering new observables is important to better understand stellar dynamics and asteroseismology.
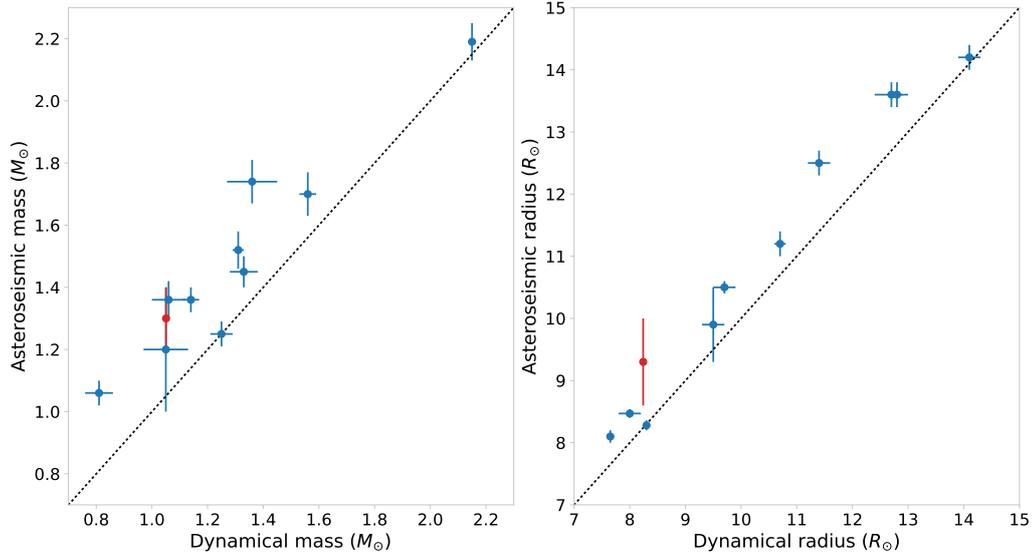
Fig. 2. Comparison between dynamical measurements and seismic inferences of masses and radii. *Left panel:* Seismic masses as a function of dynamical masses. The blue dots correspond to the sample of Gaulme et al. (2016) and the red dot, to our measurements on KIC 4054905. The grey dotted line represents the line on which $M_{\mathrm{seismo}} = M_{\mathrm{dyn}}$. *Right panel:* Seismic radii as a function of dynamical radii.

References

Abdul-Masih, M., Prša, A., Conroy, K., et al. 2016, AJ, 151, 101

Baglin, A., Auvergne, M., Boisnard, L., et al. 2006, in COSPAR Meeting, Vol. 36, 36th COSPAR Scientific Assembly

Beck, P. G., Bedding, T. R., Mosser, B., et al. 2011, Science, 332, 205

Beck, P. G., Montalban, J., Kallinger, T., et al. 2012, Nature, 481, 55

Borucki, W. J., Koch, D., Basri, G., et al. 2010, Science, 327, 977

Brown, T. M., Gilliland, R. L., Noyes, R. W., & Ramsey, L. W. 1991, ApJ, 368, 599

Corsaro, E., Mathur, S., García, R. A., et al. 2017, A&A, 605, A3

García, R. A., Hekker, S., Stello, D., et al. 2011, MNRAS, 414, L6

García, R. A., Mathur, S., Pires, S., et al. 2014, A&A, 568, A10

Gaulme, P., McKeever, J., Jackiewicz, J., et al. 2016, ApJ, 832, 121

Husser, T.-O., Wende-von Berg, S., Dreizler, S., et al. 2013, A&A, 553, A6

Kjeldsen, H. & Bedding, T. R. 1995, Astrophysical Journal, 293, 87

Mathur, S., García, R. A., Régulo, C., et al. 2010, A&A, 511, A46

Mosser, B., Barban, C., Montalbán, J., et al. 2011, A&A, 532, A86

Mosser, B., Dziembowski, W. A., Belkacem, K., et al. 2013, A&A, 559, A137

Pires, S., Mathur, S., García, R. A., et al. 2015, A&A, 574, A18

Raskin, G., van Winckel, H., Hensberge, H., et al. 2011, A&A, 526, A69

Rawls, M. L., Gaulme, P., McKeever, J., et al. 2016, ApJ, 818, 108

Rodrigues, T. S., Bossini, D., Miglio, A., et al. 2017, MNRAS, 467, 1433

Rucinski, S. M. 2002, AJ, 124, 1746

Southworth, J. 2013, A&A, 557, A119